

ABSTRACT

We examine the mechanism for the violation of the Mott-Ioffe-Regel (MIR) condition for the maximal metallic resistivity near the interaction-driven Mott transition. We find that the large resistivity appears due to the large scattering rate which destroys well-defined quasi-particles, while the Drude-like peak in the optical conductivity persists well beyond the MIR limit, similar as in the experiments on VO₂. The disorder induces effective local carrier doping which moves the system away from the Mott insulator, increasing the conductivity, in agreement with the recent experiments on X-ray irradiated organic charge-transfer salts.

Disordered Hubbard model

$$H = - \sum_{ij,\sigma} t_{ij} d_{i\sigma}^\dagger d_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow} + \varepsilon_i \sum_{i\sigma} n_{i\sigma}$$

DMFT approximation

Local Green's function $G_l(\tau - \tau') = -\langle T d_\sigma(\tau) d_\sigma^\dagger(\tau') \rangle_{S_{eff}}$ is obtained from the corresponding Anderson impurity model action [1]

$$S_{eff} = - \int_0^\beta d\tau \int_0^\beta d\tau' \sum_\sigma d_\sigma^\dagger(\tau) \mathcal{G}_{0i}^{-1}(\tau - \tau') d_\sigma(\tau') + U \int_0^\beta d\tau n_\uparrow(\tau) n_\downarrow(\tau),$$

$$\mathcal{G}_{0i}^{-1}(i\omega_n) = i\omega_n + \mu - \varepsilon_i - \Delta(i\omega_n)$$

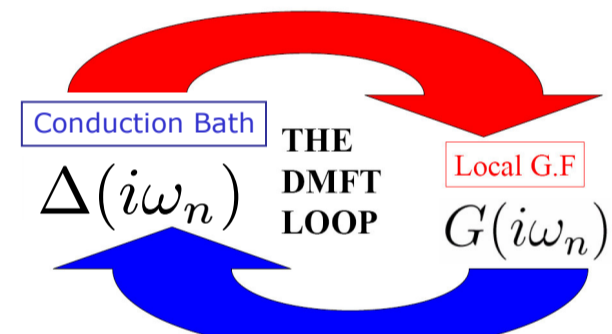
"Weiss field" random site energy conduction bath

Disorder averaging within CPA

$$G(i\omega_n) = \frac{1}{N} \sum_{i=1}^N G_i(i\omega_n)$$

DMFT self-consistency loop

EFFECTIVE LOCAL IMPURITY PROBLEM



SELF-CONSISTENCY CONDITION

$$G(i\omega_n) = \int d\varepsilon \frac{D(\varepsilon)}{\Delta(i\omega_n) + G(i\omega_n)^{-1} - \varepsilon}$$

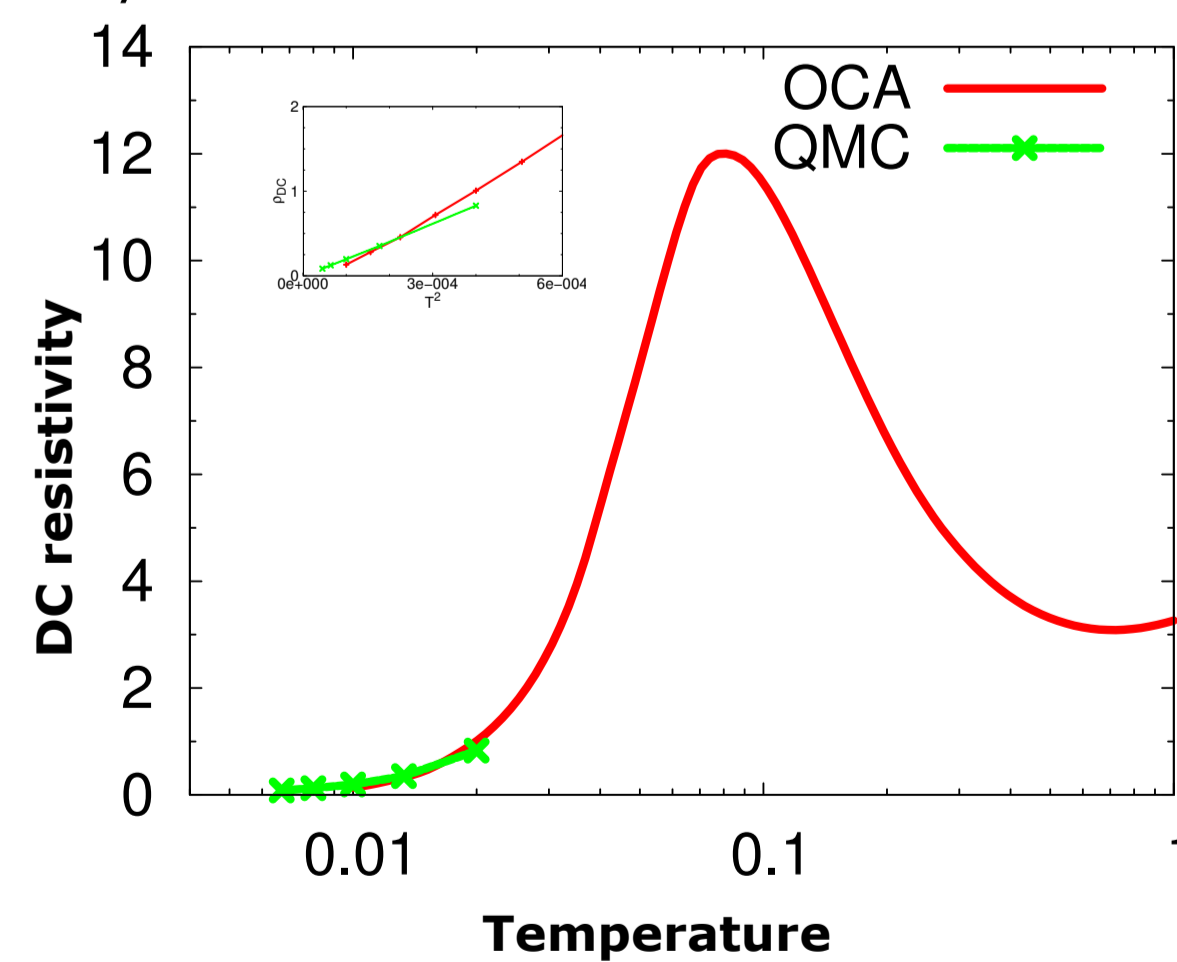
Kubo formula in DMFT approximation

$$\sigma(i\omega) = \frac{1}{\omega} \int_{-\infty}^{+\infty} d\varepsilon \int_{-\infty}^{+\infty} d\nu \int_{-\infty}^{+\infty} d\nu' D(\varepsilon) \rho(\varepsilon, \nu) \rho(\varepsilon, \nu') \frac{f(\nu) - f(\nu')}{\nu - \nu' + i\omega}$$

DC resistivity: $\rho_{DC} = \sigma^{-1}(\omega \rightarrow 0)$

Impurity solvers

Anderson impurity model is solved using One-Crossing Approximation (OCA) (generalized NCA) [2,3], and the results are cross-checked with numerically exact Continues Time Quantum Monte Carlo (CTQMC) [4] impurity solver.

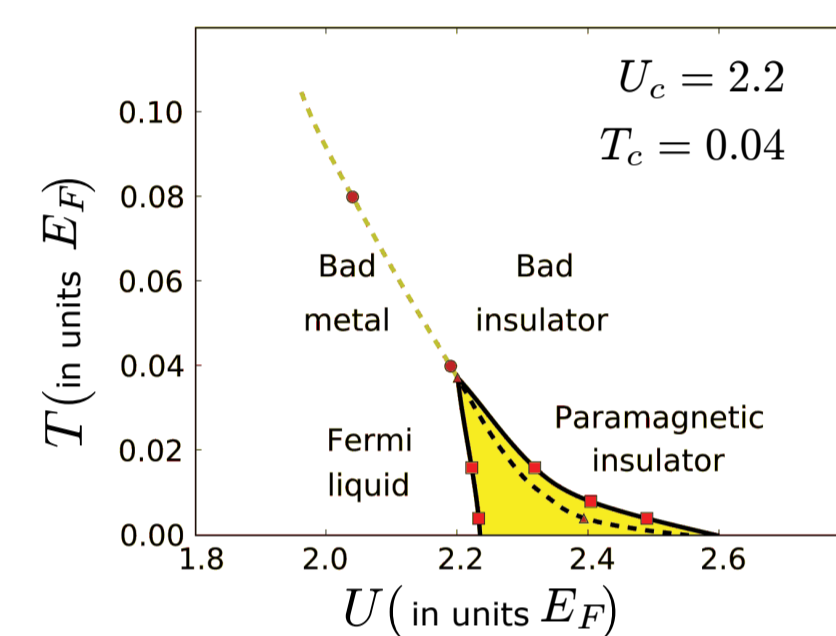


Violation of the MIR condition

• **Mott-Ioffe-Regel limit** is the maximal metallic resistivity, ρ_{MIR} , which is reached when the mean free path of the electron becomes comparable to the lattice spacing [5] $l \sim a \Leftrightarrow E_F \tau \sim 1$.

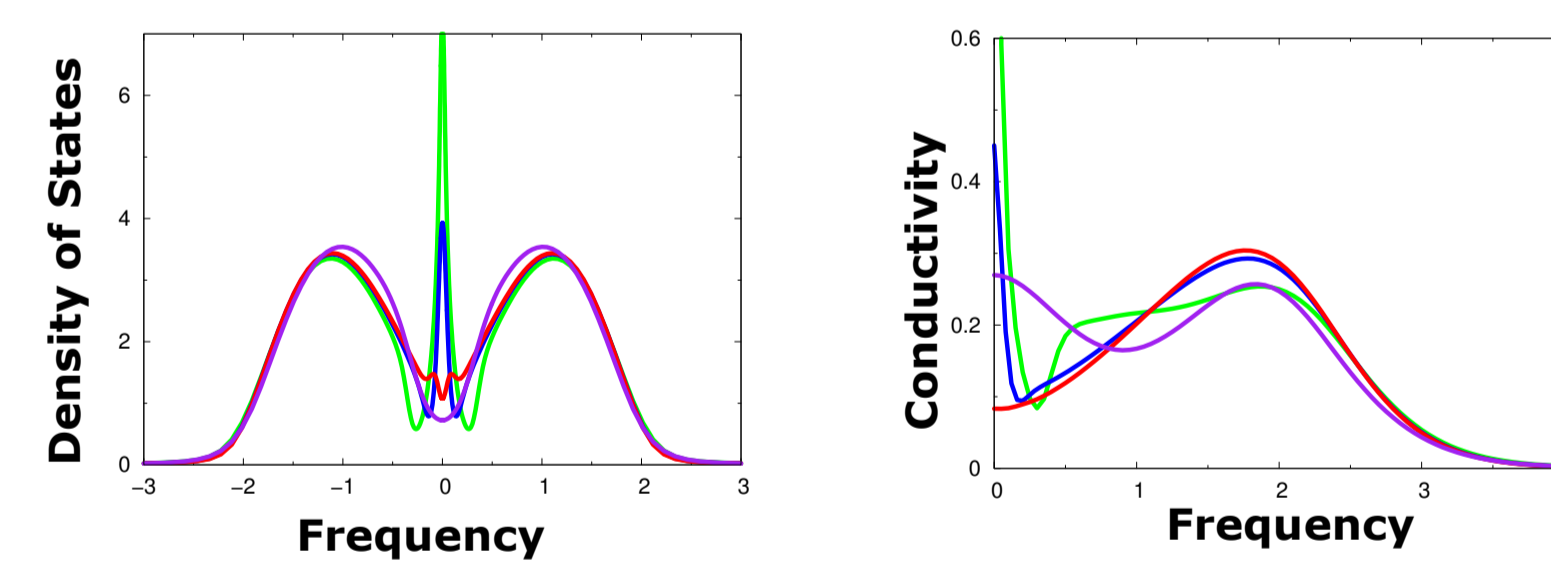
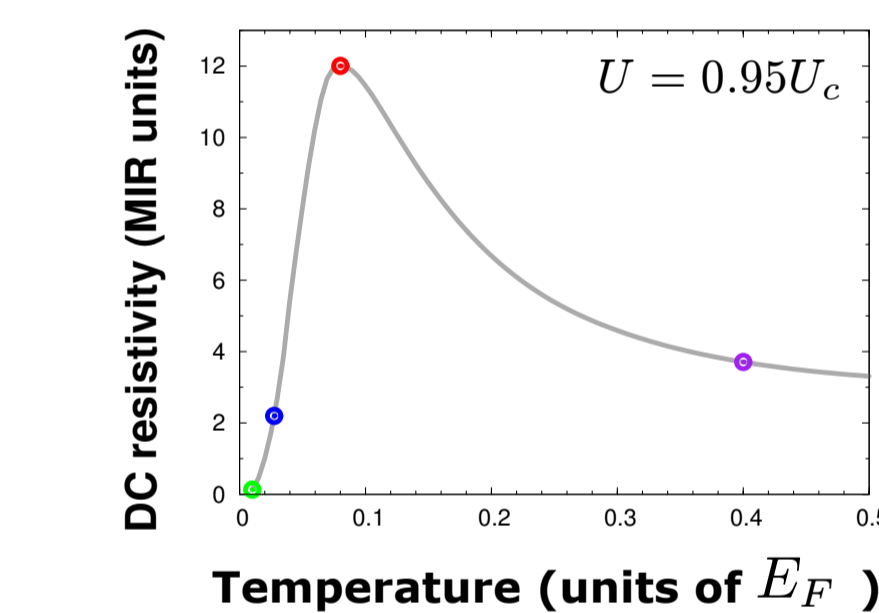
According to semi-classical arguments, resistivity cannot be larger than MIR limit.

Mott metal-insulator transition

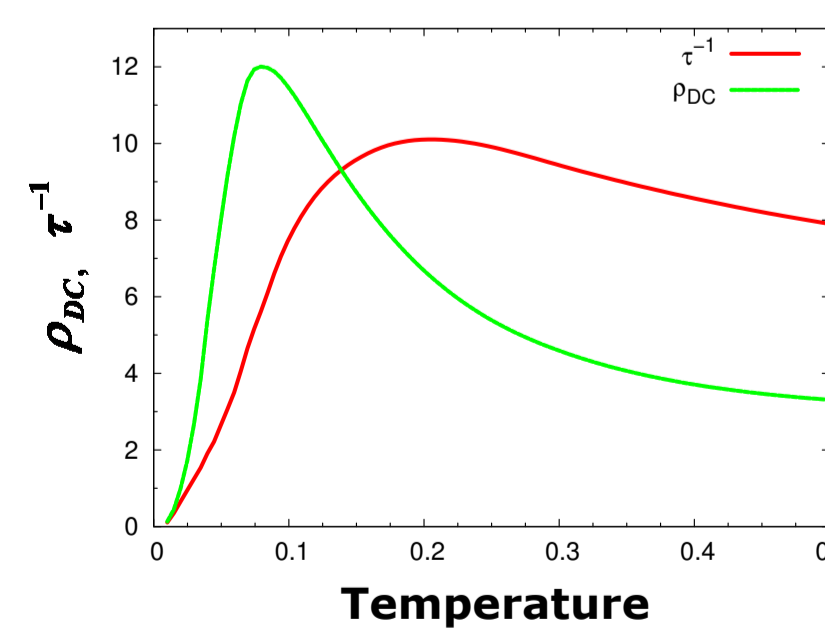


Results near the critical point

- Fermi liquid at low temperatures.
- Drude-like peak persists well above the MIR limit.
- Maximal resistivity is an order of magnitude larger than ρ_{MIR} .

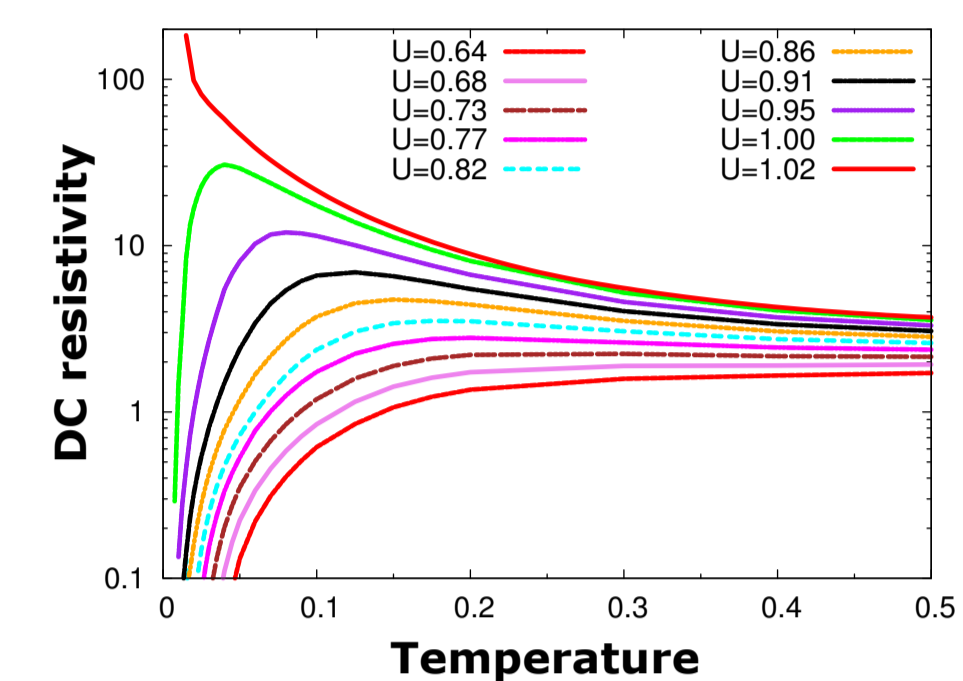


• Large **scattering rate**, $\tau^{-1} = -\frac{1}{2} \text{Im} \Sigma(0^+)$, gives the main contribution to the resistivity temperature dependence, and causes the violation of the MIR condition. Above the MIR limit, $\tau^{-1} > E_F$, and therefore, well-defined quasi-particles do not exist. Similar conclusions are obtained from experiments on VO₂ [6], and κ -organics [7].

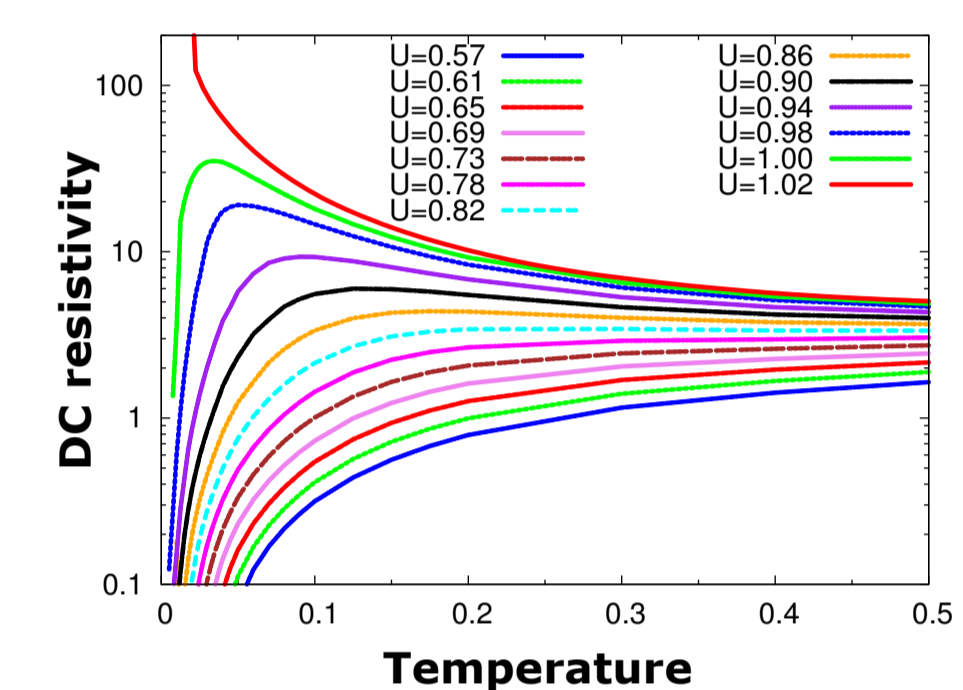


Effects of disorder

- **Disorder** effectively increases the bandwidth, and critical interaction U_c . Maximal value of ρ_{DC} , and the critical temperature do not change appreciably with disorder.



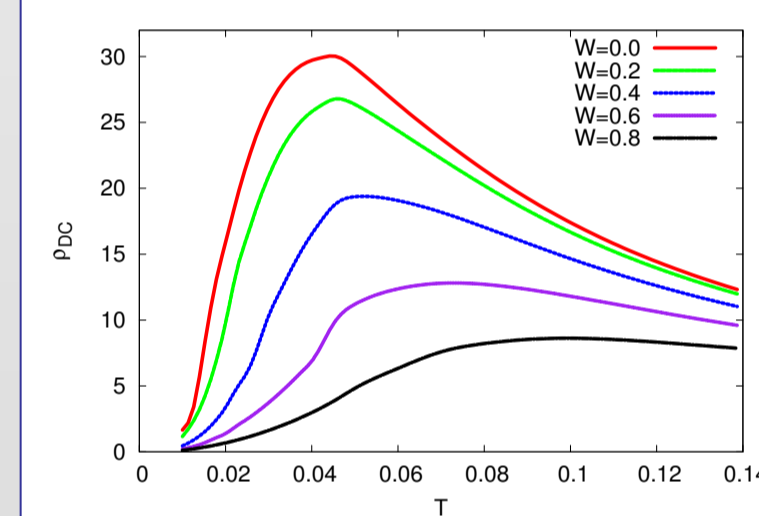
Pure case
 U in units U_c^0



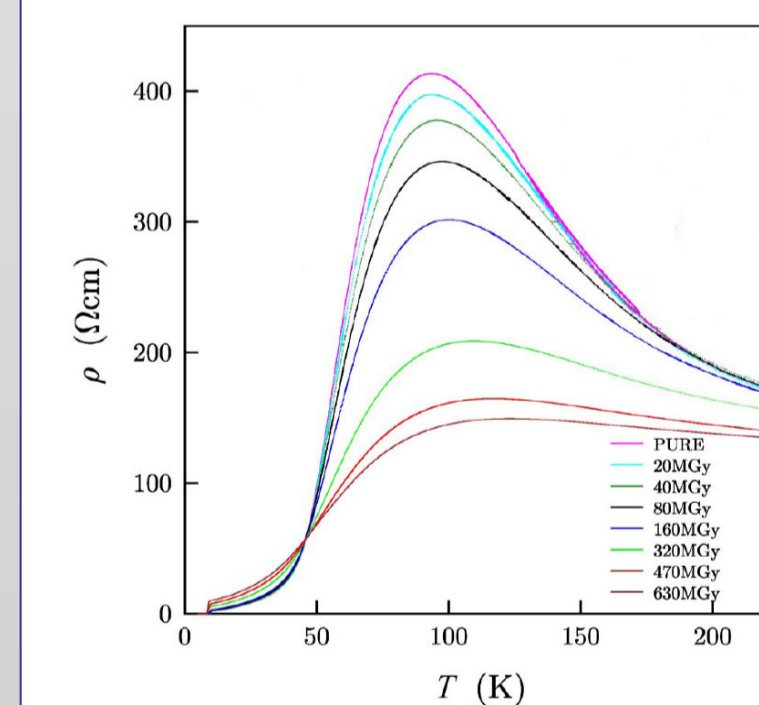
Disordered case
Uniformly distributed disorder in $(-W/2, W/2)$, U in units $U_c(W)$.

Here $U_c^0 = 2.2$, $W = 1$, $U_c = 2.45$. Increase in U_c with W agrees with earlier estimates [8].

Resistivity for fixed U



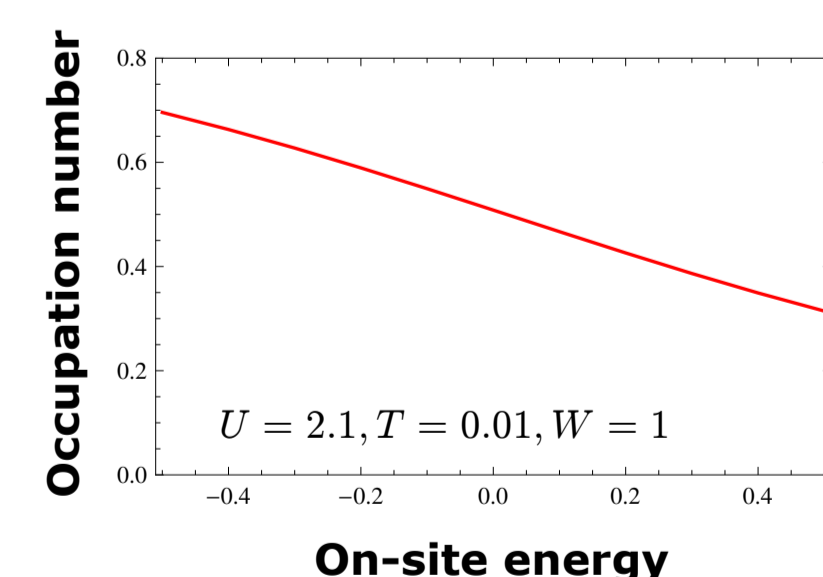
U is fixed to U_c^0 , disorder is varied



Experiments on X-ray irradiated κ -(BEDT-TTF)₂Cu(SCN)₂ [9]

Effective local carrier doping by disorder

Random site energies change the local occupation number from half-filling, reducing the correlation effects and decreasing the resistivity. Similar physical picture is proposed as the explanation of the experiments on κ -organics [10].



ACKNOWLEDGEMENTS

This work was supported in part by the Serbian Ministry of Science and Technological Development, under project No OI141035. M. Radonjić has received support under FP6 Center of Excellence grant CX-CMCS. Numerical results were obtained on the AEGIS e-Infrastructure, supported in part by FP7 projects EGEE-III and SEE-GRID-SCI. D.T. acknowledges support from the NATO Science for Peace and Security Programme Reintegration Grant No. EAP.RIG.983235.

REFERENCES

- [1] A. Georges, G. Kotliar, W. Krauth, and M. Rozenberg, Rev. Mod. Phys. **68**, 13 (1996).
- [2] Th. Pruschke and N. Grewe, Z. Phys. B: Condens. Matter **74**, 439 (1989).
- [3] K. Haule, S. Kirchner, J. Kroha, and P. Wolfe, Phys. Rev. B **64**, 155111 (2001).
- [4] K. Haule, Phys. Rev. B **75**, 155113 (2007).
- [5] O. Gunnarsson, M. Calandra, and J. E. Han, Rev. Mod. Phys. **75**, 1085 (2003).
- [6] D. N. Basov et al, Phys. Rev. B **74**, 205118 (2006).
- [7] J. Merino, et al, Phys. Rev. Lett. **100**, 086404 (2008).
- [8] D. Tanasković, et al, Phys. Rev. Lett. **91**, 0666031 (2003).
- [9] J. G. Analytis, et al, Phys. Rev. Lett. **96**, 177002 (2006).
- [10] T. Sasaki, et al, Phys. Rev. Lett. **101**, 206403 (2008).